## NumQM Spring2022

## HW3 due 01/05/2022

( +1 points if you solve the problems with Julia)
( +1 points for handing in on time)

## Q1 Dual of 3D Ising model

a) Show that the low temperature expansion of the 3D Ising model is equivalent to the high temperature expansion of a $3 \mathrm{D} \mathrm{Z}(2)$ gauge theory. (work out the first 2 correction terms)
b) What about the high temperature expansion of the 3D Ising model? Is it equivalent to the low temperature expansion of the $\mathrm{Z}(2)$ gauge theory?
c) Argue that the the dual of $4 \mathrm{D} \mathrm{Z}(2)$ gauge theory is another $\mathrm{Z}(2)$ gauge theory.

Q2
a) Show that the heat capacity

$$
c=\frac{\partial}{\partial T} \epsilon
$$

where $\epsilon=\langle E\rangle / N$, can be obtained from measuring the fluctuation of the total energy, i.e.

$$
c=\beta^{2} \frac{1}{N}\left(\left\langle E^{2}\right\rangle-\langle E\rangle^{2}\right) .
$$

b) Numerically realize the equivalence of the two for Ising Model in 4D.

Q3
a) Consider

$$
G(\lambda)=e^{\lambda H} A e^{-\lambda H},
$$

derive the differential equation

$$
\frac{d}{d \lambda} G(\lambda)=[H, G(\lambda)]
$$

b) Show that

$$
G(1)=A+[H, A]+\frac{1}{2!}[H,[H, A]]+\ldots
$$

c) Choose two $3 \times 3$ matrices A and H and use the Euler method (1st order) to solve for $G(\lambda=1)$ numerically. Compare with a direct computation of $G(1)$.
d) Prove the Baker-Campbell-Hausdorff formula:

$$
\begin{aligned}
e^{A} e^{B} & =e^{C} \\
C & =A+B+\frac{1}{2}[A, B]+\frac{1}{12}([A,[A, B]]-[B,[A, B]])+\ldots
\end{aligned}
$$

## Q4

The N-body Lorentz Invariant phase space (LISP) is defined as

$$
\begin{gathered}
\phi_{N}\left(s=P^{2}\right)=\int \frac{d^{3} p_{1}}{(2 \pi)^{3}} \frac{1}{2 E_{1}} \frac{d^{3} p_{2}}{(2 \pi)^{3}} \frac{1}{2 E_{2}} \cdots \frac{d^{3} p_{N}}{(2 \pi)^{3}} \frac{1}{2 E_{N}} \times \\
(2 \pi)^{4} \delta^{4}\left(P-\sum_{i} p_{i}\right) .
\end{gathered}
$$

Note that $E_{j}=\sqrt{p_{j}^{2}+m_{j}^{2}}$.
a) Show that the $N=2$-body phase space is given by

$$
\begin{aligned}
\phi_{2}(s) & =\frac{q(s)}{4 \pi \sqrt{s}} \\
q(s) & =\frac{1}{2} \sqrt{s} \sqrt{1-\frac{\left(m_{1}+m_{2}\right)^{2}}{s}} \sqrt{1-\frac{\left(m_{1}-m_{2}\right)^{2}}{s}}
\end{aligned}
$$

b) Numerically compute the 2-body phase space integral and check with the analytic result. Plot the results as a function of $\sqrt{s}$.

## Q5

$W=\ln Z$ is extensive.
Introducing an external field $h$ to the integral

$$
Z(h)=\int_{-\infty}^{\infty} d x e^{-V\left(x^{2}+x^{4}\right)+V h x}
$$

and define

$$
Z_{n}=\frac{1}{Z(h)} \partial_{h}^{n} Z(h)=\left\langle x^{n}\right\rangle \times V^{n}
$$

and

$$
\begin{aligned}
W(h) & =\ln Z(h) \\
W_{n} & =\partial_{h}^{n} W(h) .
\end{aligned}
$$

a) Show that (see Sem01 HW02 Q4, but this time we do not set $h \rightarrow 0$ at the end.)

$$
\begin{aligned}
& W_{1}=Z_{1} \\
& W_{2}=Z_{2}-Z_{1}^{2} \\
& W_{3}=Z_{3}-3 Z_{1} Z_{2}+2 Z_{1}^{3} \\
& W_{4}=Z_{4}-4 Z_{1} Z_{3}-3 Z_{2}^{2}+12 Z_{1}^{2} Z_{2}-6 Z_{1}^{4}
\end{aligned}
$$

b) Verify these relations numerically, i.e. taking numerical derivatives on $W(h)$ and performing corresponding numerical integrations

$$
Z_{n}=\frac{1}{Z(h)} \int_{-\infty}^{\infty} d x V^{n} x^{n} e^{-V\left(x^{2}+x^{4}\right)+V h x}
$$

Plot them as functions of $h$ in the range of $-2: 2$. (Take $\mathrm{V}=2$.)
c) Verify the linked cluster theorem: $W_{n} \propto V$ at large $V$, i.e. $W_{n} / V=w_{n}$ becomes intensive.

